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# About 250/285 GHz push–push oscillator using differential gate equalisation in digital 65-nm CMOS

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**Abstract:** This study presents a push–push oscillator architecture based on differential gate equalisation to enhance the oscillation frequency while providing relatively high output power with ultra-compact layout form factor. The frequency enhancement is derived as a function of the equivalent RLC model of the oscillator's main constituents. The proposed principle is applied to a terahertz oscillator in the 200–300 GHz range to mitigate the excessive substrate and skin effect losses in standard digital 65-nm complementary metal–oxide–semiconductor technology at such high frequencies. The design concept is validated using two single-stage push–push oscillators. The first oscillator shows –8.1 dBm output power at 250 GHz oscillation frequency and –106.8 dBc/Hz phase noise at 10 MHz offset while consuming 76 mW power from 1.5 V DC supply voltage. The chip area is  $200 \times 250 \mu\text{m}^2$ . The second oscillator provides –14.8 dBm output power at 285 GHz and –106 dBc/Hz phase noise at 10 MHz offset with 80 mW power consumption from 1.5 V DC supply. The chip area is  $200 \times 200 \mu\text{m}^2$ .

## 1 Introduction

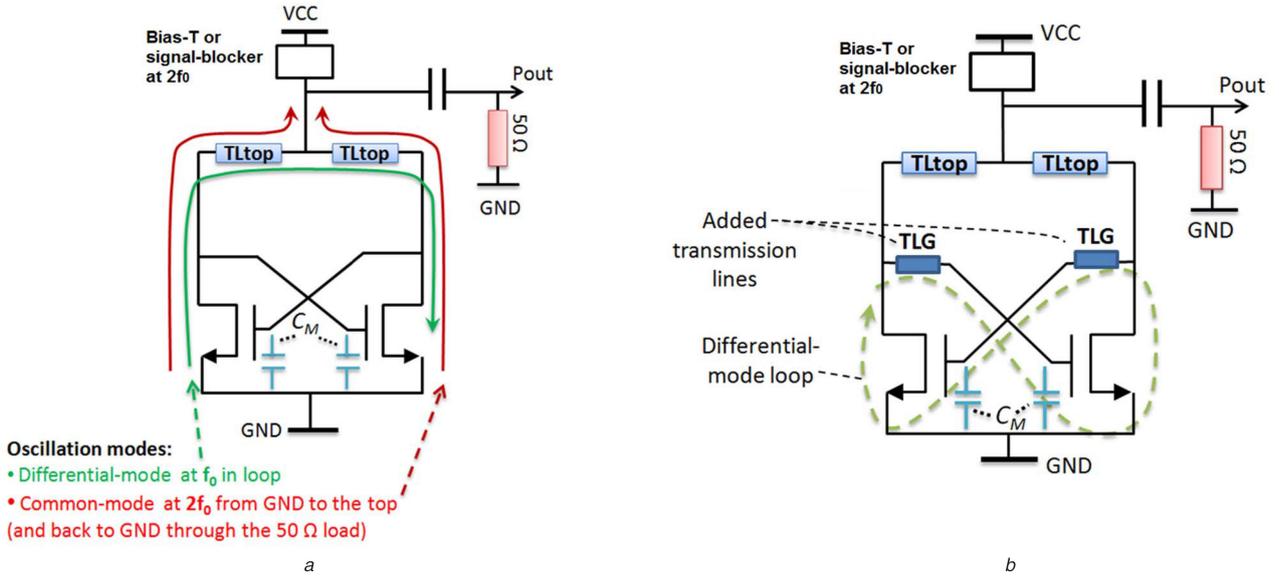
The unique properties for the electromagnetic (EM) waves in the terahertz and sub-mmWave frequency ranges (penetration, absorption, scattering etc.) promise disruptive technological advances in short-distance high data-rate communication, spectroscopy, bio-medical diagnosis, imaging and others. Existing THz systems are mostly bulky, expensive and based on photonic realisations [1]. Integrated systems using low-cost standard integrated circuit (IC) technologies have become a requirement to provide compact, portable, cost-effective and adaptive commercial solutions [2]. It remains that most of the used standard IC processes [complementary metal–oxide–semiconductor (CMOS) or bipolar] have their own frequency limitations and increased parasitic losses at such high frequencies, which seriously reduce the output transmission power and impede the THz solutions from moving towards low-cost commercial products.

One approach of generating power at THz frequencies is based on amplifier–multiplier chains with the input taken from a stable low-frequency source to generate the output signal by mixing effects at higher harmonic frequencies as in [3–6]. This approach suffers from high DC power and large chip area. In addition, the phase noise of the output is a function of the input signal's phase noise and thus a stable input frequency source is required. Another approach for THz signal generation is direct generation using high-frequency coupled oscillators to extract the higher order harmonics of the fundamental oscillation frequency. A small chip size could be achieved with low DC power and the signal source in this case is self-contained. Push–push [7–9] and self-feeding harmonic [10] oscillators use balanced structures and symmetric layouts to extract the second harmonic. Triple-push [11–13] and quad-push [14] oscillator topologies have also been utilised to achieve higher oscillation frequencies at the third and fourth harmonics, respectively. However, triple-push oscillators suffer from complexity in generating a symmetric layout while the output power of the fourth harmonic in the case of quad-push oscillators is relatively small. To increase the output power, either multiple core oscillators are coupled [13] or arranged in an array format [15].

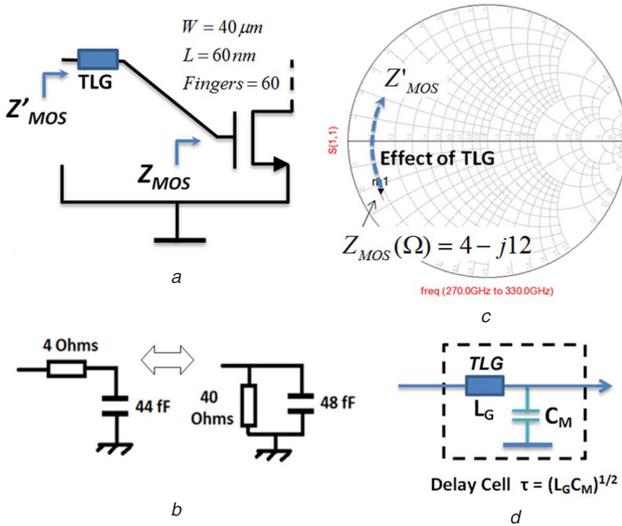
In THz and sub-mmWave frequency ranges, substrate losses and skin effects decrease the quality factor of passive devices and varactors making it more challenging to extract the highest power from the transistors within the oscillator while targeting a small layout form factor. This is particularly true for digital CMOS processes which tend to have a thin back end of line. Design examples proposed in the literature include using a slow-wave coplanar waveguide (CPW) which has lower loss than a standard or grounded CPW to implement the tank [13]. Varactor-based techniques are applied using hyper-abrupt junction varactor to increase the linearity of  $C$ – $V$  characteristics or the parasitic capacitance of metal–oxide–semiconductor (MOS) transistor in [16, 17], respectively, with both oscillators operated in the 100 GHz band.

In this paper, a push–push oscillator architecture based on a differential structure is demonstrated. The proposed approach is capable of achieving a compact form factor through a simplified routing with less parasitics in addition to reduced DC power consumption. We show that the use of differential equalisation between the cross-coupled metal–oxide–semiconductor field-effect transistor gates reduces the equivalent capacitance at the gate level which results in increasing the frequency of oscillation. A frequency enhancement factor (FEF) is derived as a function of the equivalent RLC model of the oscillator design parts. The design is validated using two single-stage push–push oscillators implemented in 65 nm standard digital CMOS technology. The first oscillator shows –8.1 dBm output power at 250 GHz oscillation frequency while consuming 76 mW power from 1.5 V DC supply with  $200 \times 250 \mu\text{m}^2$  chip area. The second oscillator gives –14.8 dBm output power at 285 GHz with 80 mW power from 1.5 V DC supply and  $200 \times 200 \mu\text{m}^2$  chip area.

The paper is organised as follows. In Section 2, the general push–push design principle with differential gate equalisation is presented. In Section 3, the implementation and the different oscillator parts modelling are detailed. Section 4 gives the experimental setup description and the measurement results. Finally, conclusions are drawn in Section 5.



**Fig. 1** Push–push oscillator  
 (a) With a traditional cross-coupled active gm-core,  
 (b) With added TLG



**Fig. 2** MOS characteristics shown through  
 (a) The gate impedance at 300 GHz represented by  $Z_{MOS}$  and  $Z'_{MOS}$  after TLG,  
 (b) Equivalent  $Z_{MOS}$  model with series and parallel elements,  
 (c) Effect of the added transmission line (TLG) on transforming the equivalent MOS capacitive impedance ( $Z_{MOS}$ ) to an inductive impedance ( $Z'_{MOS}$ ).  
 (d) Delay cell formed by TLG and the equivalent MOS capacitance ( $C_M$ )

## 2 Design principle

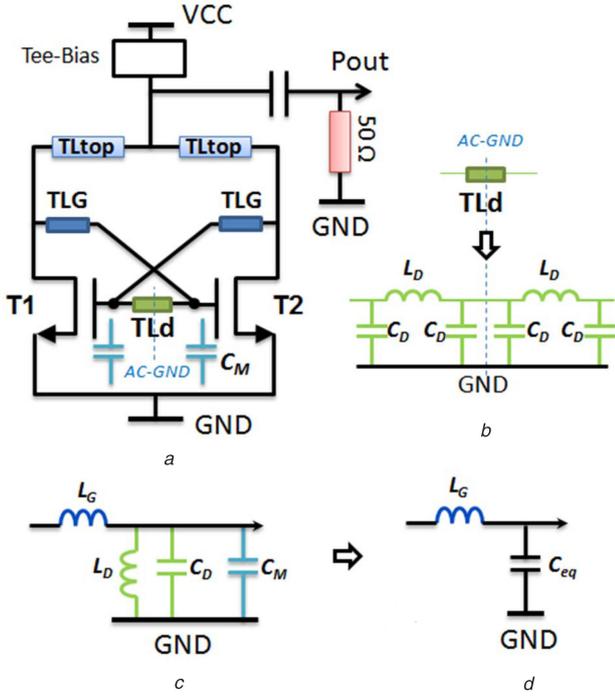
A typical oscillator architecture is based on the use of a resonating part and a transconductance gm-core providing the negative resistance to compensate for the tank losses. Fig. 1a shows a differential oscillation mode at a frequency  $f_0$ , where the signal goes through an internal loop without the need to be referenced to the ground (GND). The push–push behaviour can be described as the circulation of the common-mode (CM) signal at a frequency  $2f_0$  (or generally at even harmonics of  $f_0$ ) through a loop going from GND to the output load and back to GND. The output signal could be either electrical or radiated (output terminated with an antenna). In the case of an electrical signal, a bias-T or a signal blocker at  $2f_0$  could be used to provide the DC supply voltage (VCC). Transmission lines (TL) use at the intended frequency range is preferred as they present much lower losses than lumped elements, inductor and capacitor-based tanks [18].

To increase the output power from the oscillator, the conversion from differential mode to CM signal should be optimised. Since

part of the CM power at even harmonics (mainly at  $2f_0$ ) travelling through the gm-core from GND to the top load will be absorbed by the gm-core transistors' gates, this reduces the power delivered to the output. Adding a 'signal blocking' impedance or transmission line before the gates can prevent such losses if the realised impedance transformation goes from capacitive to inductive. This can also be expressed as a function of gate-blocking TL (TLG) parameters, i.e. the characteristic impedance  $Z_{0,TLG}$  and the electrical length  $\varphi_{TLG}$  to maximise the CM power transfer to the output at  $2f_0$ , as was also reported in [10]. Different ( $Z_{0,TLG}$ ,  $\varphi_{TLG}$ ) can be determined through simulations to satisfy the maximum output power condition. In this paper, the goal is to build on this understanding and also realise a compact layout with short TLG to minimise the associated losses with the series resistance. On the top side, the TLtop transmission lines depicted in Fig. 1b also contribute to the differential-mode circulation. However, they do not influence the frequency enhancement effect which will be introduced later. Further considerations for TLtop sizing are presented in the implementation section.

Based on the previous discussion, a TLG is added to the cross-coupled oscillator structure, as shown in Fig. 1b. The used MOS transistor (MOST) size is selected to get maximum transition frequency ( $f_T$ ) performance. With  $f_T = g_m/2\pi C_{in}$ , where  $g_m$  is the MOS transconductance and  $C_{in}$  is its input gate capacitance,  $f_T$  allows selecting optimum MOST size and bias point settings regardless of the input/output loading conditions. The MOS size was accordingly set to a width of 40  $\mu\text{m}$ , a gate length of 60 nm and folded into 60 fingers. The gate impedance characterisation, represented by  $Z_{MOS}$  in Figs. 2a–c, gives  $Z_{MOS} = 4 - j12$  at 300 GHz, which can be modelled by an equivalent capacitance  $C_M \approx 48$  fF. The effect of TLG can hardly render infinite input impedance from the gate level. However, making the transformed impedance ( $Z'_{MOS}$ ) inductive would relax the CM impedance matching. Setting TLG to an electrical length of roughly  $30^\circ$  and 50- $\Omega$  characteristic impedance is equivalent to a series inductance of  $L_G \approx 15$  pH. Considering either TLG or  $L_G$  would sensitively give the same impedance rotation starting from  $Z_{MOS}$  on the Smith chart shown in Fig. 2c. This also means that the parasitic capacitance for TLG, i.e.  $C_G \approx 6$  fF (calculated with  $Z_c = (L_G/C_G)^{1/2} = 50 \Omega$ ) can be neglected with respect to  $C_M$ . The path through TLG to the MOST can thus be considered as a delay cell formed basically by  $L_G$  and  $C_M$ .

The oscillation frequency ( $f_0$ ) is dictated by two parallel paths. The first path is through TLtop and the second one is through the



**Fig. 3** Oscillator design represented

- (a) Oscillator design with added TLd,  
(b) Equivalent dual-section  $\pi$ -model for TLd assumed lossless,  
(c) New delay cell with added elements  $L_D$  and  $C_D$  representing a half-length of TLd with respect to the middle AC-ground,  
(d) Delay cell sketched with the equivalent capacitance  $C_{eq}$  of ( $L_D$ ,  $C_D$ ,  $C_M$ )

delay cell formed by ( $L_G$ ,  $C_M$ ). It can be demonstrated (analytically and by simulation) that the resulting oscillation frequency is given by

$$f_0 = \left( \frac{f_1}{a_1} // \frac{f_2}{a_2} \right) \quad (1)$$

where  $a_1$  and  $a_2$  are coefficients that can be related to the relative impedances of TLtop and the ( $L_G$ ,  $C_M$ ) delay cell, respectively, and  $f_1$  and  $f_2$  are the relative frequencies set by the delays through these TLLs, such that

$$f_1 = \frac{1}{2} \frac{1}{2\pi\sqrt{L_G C_M}} \quad (2)$$

and

$$f_2 = \frac{1}{4} \frac{v}{2\pi l_{top}} \quad (3)$$

where  $l_{top}$  is the physical length of TLtop and  $v$  is the light velocity in the considered material.

In this paper, the introduced frequency enhancement effect is relative to the oscillation frequency increase through  $f_1$ . Thus, the following developments will be only in regard to this frequency. Note that the calculated frequency (2) does not take into account the equivalent capacitance at the MOS drain, which would even result in a lower fundamental frequency. As mentioned earlier, opting for longer TLG to present higher impedance or ideally an open circuit at the gate would also increase  $L_G$  and lower  $f_0$ . This behaviour defines a trade-off between maximising the output power and achieving a high oscillation frequency.

A similar design approach has been demonstrated in [9, 10, 14]. In [9], the gate blocking is applied on eight single-ended stages to extract the second harmonic through direct power combining. In [10], the gate blocking is employed in eight differential oscillators each generating the second harmonic along with radiated power combining. In [14], the gate-blocking impedance is used in a

differential structure and connected in an oscillator loop of four cells. In this case, the fourth harmonic ( $4f_0$ ) of the oscillation frequency is exploited to generate the output signal (quadruple push oscillator). In all these cases, having a large number of stages gives higher routing complexity, more parasitics, higher DC power consumption and an increasing chip area.

The proposed oscillator design is based on a differential structure to achieve compact form factor, simplified routing with less parasitics and reduced DC power consumption. A differential transmission line (TLd) is added between the cross-coupled transistor gates. Its equivalent series inductance ( $L_D$ ) cancels part of the gate capacitance ( $C_M$ ) in addition to the TLd's own distributed parasitic capacitance ( $C_D$ ), as shown in Figs. 3a and b. This effect could be seen as a differential equalisation where by reducing the equivalent capacitance at the gate level, a higher oscillation frequency is obtained. Note that TLd has a negligible effect on the common mode behaviour as long as its own parasitic capacitance ( $C_D$ ) can be made negligible compared to the equivalent gate capacitance ( $C_M$ ).

The transmission line (TLd), assumed to be lossless, can be represented by a double section  $C_D L_D C_D$   $\pi$ -model, as depicted in Fig. 3b. The modified delay cell given in Fig. 3c includes  $L_D$  and  $C_D$  elements in parallel to  $C_M$ , and refers to a half-length of TLd with respect to the middle AC-ground. The parallel combination of  $L_D$ ,  $C_D$ , and  $C_M$  is represented by an equivalent capacitance  $C_{eq}$  in the delay cell of Fig. 3d, such that

$$C_{eq} = C_M + C_D - \frac{1}{4\pi^2 f_1^2 L_D} \quad (4)$$

From (2), the new frequency ( $f'_1$ ) could be written as

$$\begin{aligned} f'_1 &= \frac{1}{2} \frac{1}{2\pi\sqrt{L_G C_{eq}}} \\ &= \frac{1}{2} \frac{1}{2\pi\sqrt{L_G (C_M + C_D - (1/4\pi^2 f_1^2 L_D))}} \end{aligned} \quad (5)$$

Rearranging (5), we get

$$f'_1 = \sqrt{\frac{(1 + 4L_G/L_D)}{(1 + C_D/C_M)}} f_1 = \text{FEF} \cdot f_1 \quad (6)$$

where  $f_1$  is the initial frequency calculated in (2). From (6), it can be inferred that a FEF between  $f'_1$  and  $f_1$  has resulted from the proposed equalisation method. For a given TLG design, the MOST size, for a higher FEF would require lowering both  $L_D$  and  $C_D$  which leads to a relatively short TLd. Sizing TLd width is a compromise between several factors. For example, finding the minimum series resistance (giving wide TLd), maximising FEF, and setting the right characteristic impedance. Thus, an acceptable value for TLd could be determined through simulations.

To illustrate the potential of the proposed technique, Fig. 4 gives the evolution contours of FEF as a function of  $L_G/L_D$  and  $C_D/C_M$  ratios. As  $C_D$  can be considered negligible or lower than  $C_M$ , this lets us consider further the left side in Fig. 4. A large range of  $L_G/L_D$  settings can thus be chosen to easily obtain a FEF larger than one, and this is mainly due to the factor 4 in the numerator in (6). If we assume, for instance, that TLd has similar characteristics to the previously considered TLG ( $L_G \approx 15$  pH,  $C_G \approx 6$  fF), the resulting FEF gets roughly equal to two.

### 3 Implementation and modelling

The oscillator is implemented in a standard eight-metals 65 nm digital CMOS technology from ST Microelectronics. The top metal height is about only  $5.7 \mu\text{m}$  from the substrate, which means that the parasitic capacitances (estimated to  $\sim 20$  aF/ $\mu\text{m}^2$  between M8 and M1) is quite high at the target frequency range. The oscillator design uses TL exclusively to avoid the power loss associated with

degraded quality factors of passive elements. Fig. 5a gives a descriptive schematic of the oscillator design and Fig. 5b shows a layer-level schematic illustrating TLG and TLd layout and connections to the MOSTs ( $T_1, T_2$ ).

The modelling of the different oscillator circuit components is critical to accurately predict the oscillation frequency and the output power. This might be challenging as the routing of the MOSTs and TLs has many superposed layers and shared zones to isolate distributed and close mutual capacitances or inductances [19].

The MOSTs are sized to 40  $\mu\text{m}$  width, 60 nm length and folded into 60 fingers. The substrate ring in M1 (part of the default library layout) is removed on one side and continued only with active (P+) layer to allow direct connection (with M1 and M2) from the gate to both TLG and TLd. This is done to achieve a minimum series resistance at the cost of slightly increased additional capacitance to the bulk. The source connection is made through parallel M2 and M3 layers to the circuit ground (GND) as shown in Fig. 5b. The capacitive parasitics due to the MOSTs' routing is extracted using

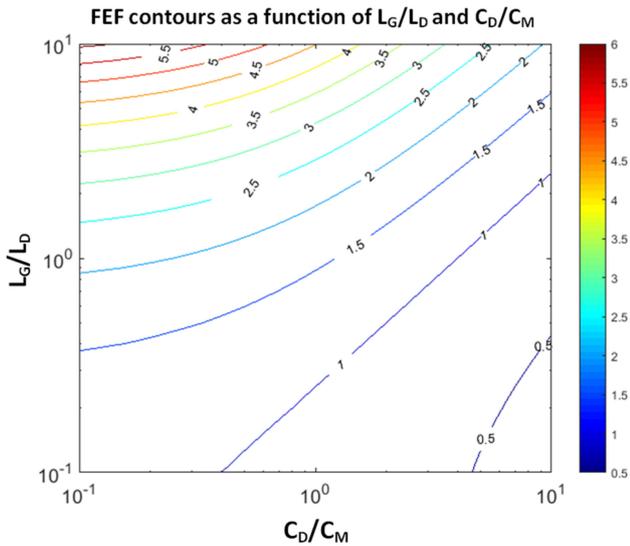


Fig. 4 FEF contours as a function of  $L_G/L_D$  and  $C_D/C_M$  ratios varied from 0.1 to 10 each

Calibre software. The overlapping vias and metal-layer series resistance connected to either the drain or the source is simulated with Momentum software to take into account the resistance increase by the skin effect. The skin depth is estimated to be approximately 0.13  $\mu\text{m}$  at 300 GHz. The extracted RC elements for each routed MOST are represented in Fig. 6a. To provide a minimum ohmic ground to the gm-core despite the very tight skin depth, two wide planes in M1 and M2 are connected over the complete circuit width ( $\sim 200 \mu\text{m}$ ) from the GND pads to the gm-core cell.

The TLG design is based on a compact and direct connection between the respective MOSTs drains and gates to minimise the inner spacing between transistors to around 7  $\mu\text{m}$ . This avoids the CM power loss by asymmetrical connections caused by lines crossing. The spacing between both TLG microstrips is set to less than 4  $\mu\text{m}$  and each trace is 3  $\mu\text{m}$  wide (Fig. 5b). The TLG layout is made with stacked metals, i.e. a M2/M3 stack starting from the gate to reach the opposite drain with a M4/M7 stack to minimise the series losses with reduced traces dimensions and form factor. The lumped RLC model for the dual TLG lines is shown in Fig. 6b. This model is extracted using S-parameter curves fitting with a four-port Momentum layout. The model is split into dual sections for improved accuracy, with  $k_f$  representing the coupling factor between the respective line inductances [20].

The TLd is placed on the bottom side of the gm-core to provide a direct access to the MOS gates in M1/M2 layers as shown in Fig. 5b. The TL is laid out in M8 layer over M1 GND and sized to 60  $\mu\text{m}$  total length and 8  $\mu\text{m}$  width. The relative RLC model (see Fig. 6c) is extracted using S-parameter curves fitting with a two-port Momentum layout. The TLd line presents a characteristic impedance of 44  $\Omega$ .

While the RLC modelling helps define the initial design, EM simulation (done with Momentum) remains unavoidable to take into account the mutual capacitances and magnetic coupling between the different design components. The EM simulation should be extensive and includes all elements without compromising the meshing accuracy or increasing excessively the simulation time. In this design, EM simulations are done on two main parts: the *first* comprises the metal connections above the MOSTs, TLG, TLG, along with the surrounding ground plane as shown in Fig. 6d, and the *second* regards the top TLs (TLtop) connection to the ground-signal-ground measurement pads as shown in Fig. 7a.

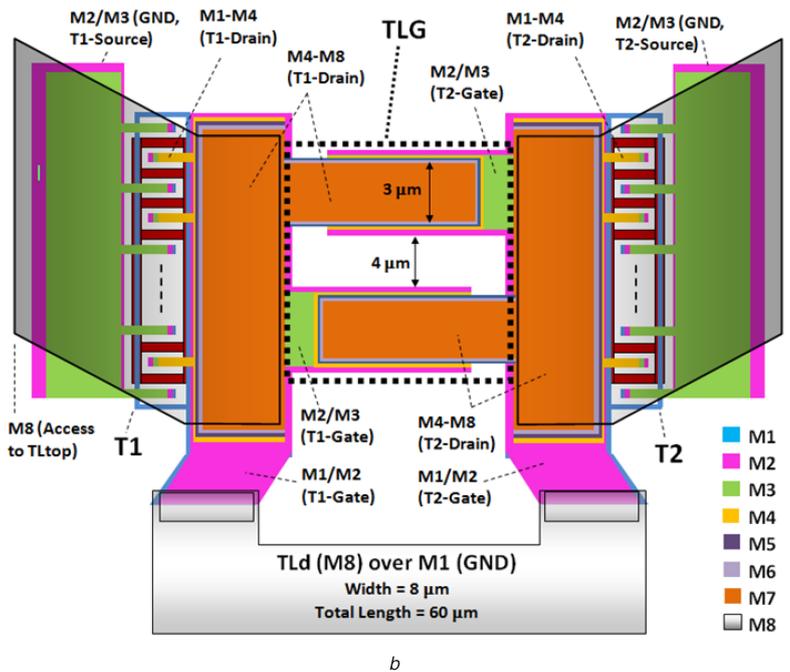
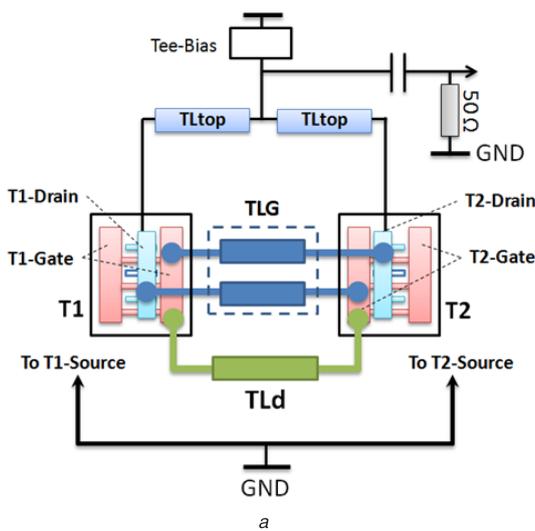
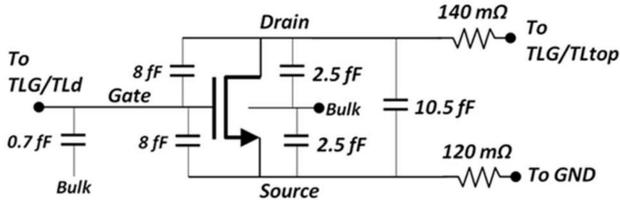
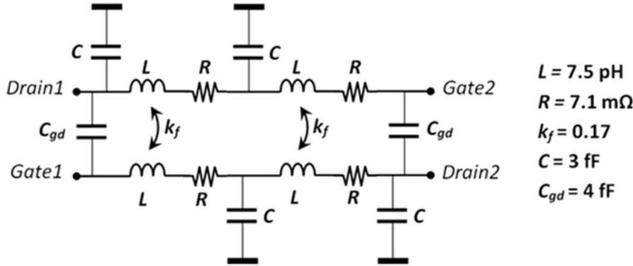


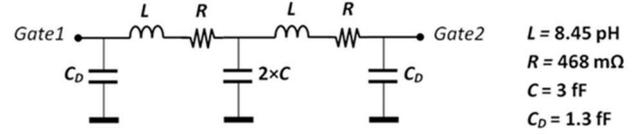
Fig. 5 Descriptive schematic of the oscillator design showing (a) TLd and TLG connections to the MOSTs, (b) Layer-level layout of the different connections and TLs (scale not representative of the real dimensions)



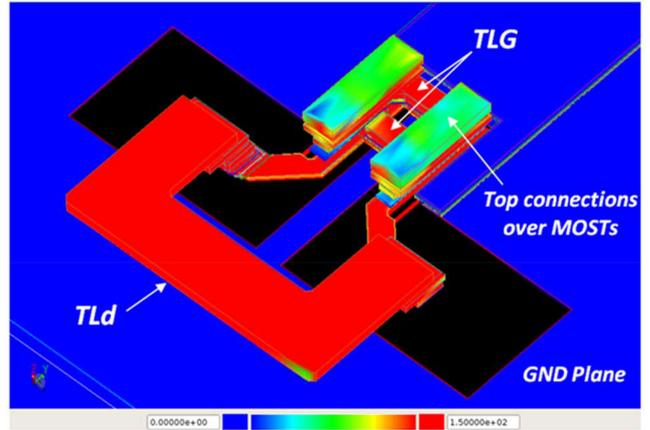
a



b



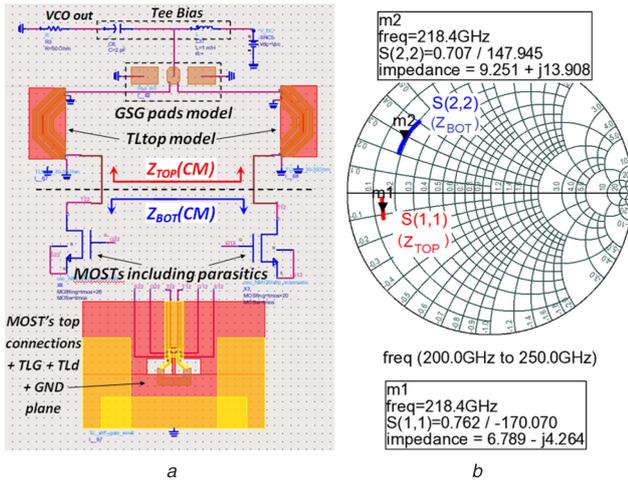
c



d

**Fig. 6** Extracted lumped RLC model

- (a) MOST routing,
- (b) Dual-section model for TLG,
- (c) TLd,
- (d) EM-simulated Momentum layout including TLd, TLG, the top MOSTs' connections, and the surrounding GND plane



**Fig. 7** ADS simulation showing

- (a) The used schematic with the different EM-simulated-based models,
- (b) Partial CM matching between the bottom oscillator impedance ( $Z_{BOT}$ ) and the top part impedance ( $Z_{TOP}$ ) at 218 GHz frequency

The TLtop model was included as two separate parts in the oscillator schematic (Fig. 7a) due to negligible mutual coupling between both lines with grounded sidewalls and bottom-M1 waveguide structure [21]. This also reduces the simulation time needed for optimisation due to the large layout area that would have taken a single EM-model for a dual TLtop layout. Each TLtop is sketched with a 12  $\mu\text{m}$  wide M8 layer over 165  $\mu\text{m}$  long track from the MOST to the signal pad. Shielding sidewalls are formed with M8-to-M1 metal stack and set to 13  $\mu\text{m}$  width and 6  $\mu\text{m}$  spacing away from the signal track. The TL gives around 40  $\Omega$  characteristic impedance.

Starting with  $\lambda/2$  length at  $2f_0$  for TLtop lines provides 50- $\Omega$  CM impedance for  $Z_{TOP}$  (looking upwards in Fig. 7a). With an inductive CM impedance presented by the entire bottom part of the oscillator ( $Z_{BOT}$  in Fig. 7b), a matching issue comes up with a non-optimum CM power transfer situation. A solution would consist of shortening TLtop length to allow  $Z_{TOP}$  rotation on the Smith chart

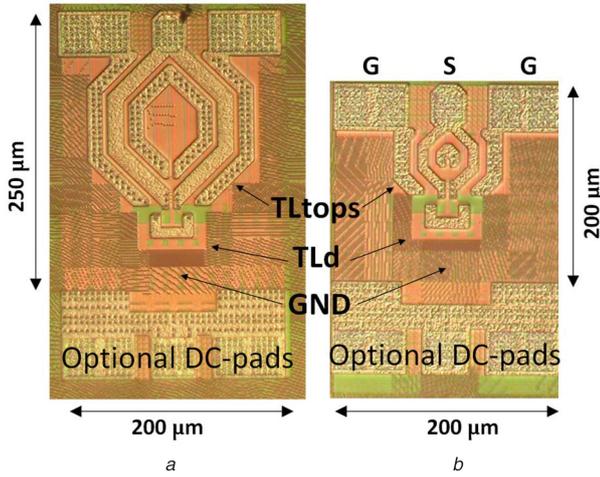
(see Fig. 7b) in order to place it on the conjugate side of  $Z_{BOT}$  for optimum CM matching. The counterpart of this design change is the partial contribution of TLtop to the oscillation frequency ( $f_0$ ) by creating a parallel differential path to the gm-core's one. Managing both constraints in terms of CM impedance matching and oscillation frequency sets TLtop length to slightly less than  $\lambda/4$  (or  $\sim 165 \mu\text{m}$ ) for a final simulated frequency at 218 GHz and  $-6.1$  dBm output power under 1.5 V DC supply voltage. The partial CM matching in this case is shown in Fig. 7b. This oscillator version is termed (OSC1). Another version (OSC2) of this oscillator design has a shortened TLtop length of 70  $\mu\text{m}$  to assess the effect on the frequency shift. TLtop is modified without any matching optimisation effort and with an expected power performance penalty. The simulated oscillation frequency for OSC2 showed around 250 GHz with about  $-9$  dBm output power from 1.5 V DC supply voltage. The main goal of this VCO design is to illustrate the fundamental frequency enhancement effect while targeting a very small layout form factor. Additional effort is needed to increase the output power via on-chip or spatial combining.

## 4 Measurement results

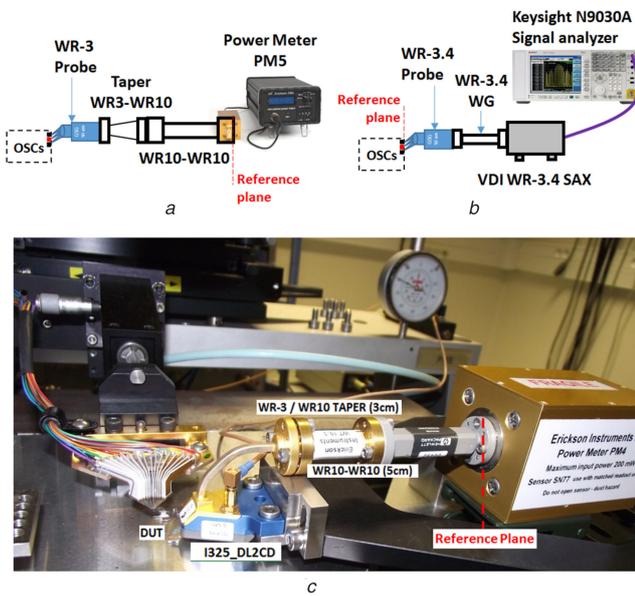
Micrographs of the implemented oscillators OSC1 and OSC2 are shown in Figs. 8a and b, respectively. Both oscillators were directly probed on chip. The effective area is 200  $\mu\text{m} \times 250 \mu\text{m}$  and 200  $\mu\text{m} \times 200 \mu\text{m}$  for OCS1 and OSC2, respectively.

The testing setups are given in Figs. 9a and b for power and frequency measurements, respectively. Erickson PM5 power meter from Virginia Diodes Inc. (VDI) was used for power measurements, while VDI WR-3.4 SAX down-converter in addition to Keysight N9030A signal analyser was employed for frequency and phase-noise measurements. Both oscillators were tested in free-running mode without any frequency locking.

Measured oscillation frequencies for both oscillators showed an upward shift compared to simulations. OSC1 has an oscillation frequency from 248 to 251 GHz for a supply voltage (VDD) varied from 1.2 to 1.5 V as shown in Fig. 10a. This is compared to 218 GHz from simulations. The deviation in frequency can be due to conservative modelling (mainly with the MOST parasitics) and to the difficulty of separating shared or overlapping connections as well as too close mutual capacitances or magnetic coupling that would have been accounted multiple times. OSC2 presents about



**Fig. 8** Implemented oscillator micrographs  
(a) OSC1,  
(b) OSC2

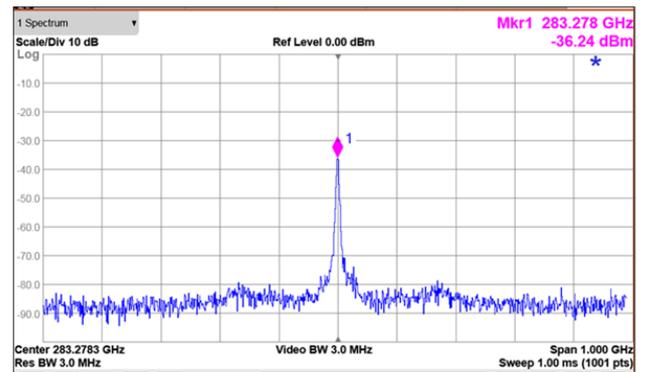
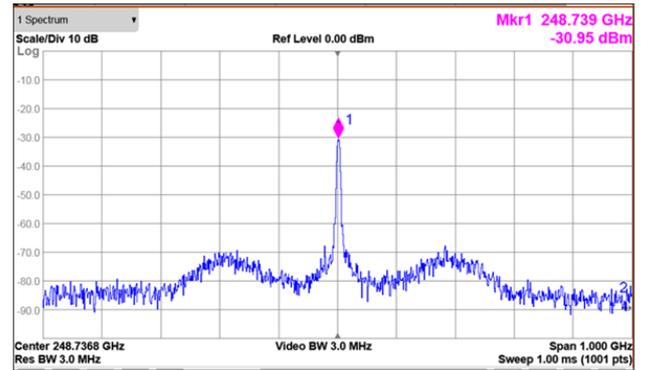
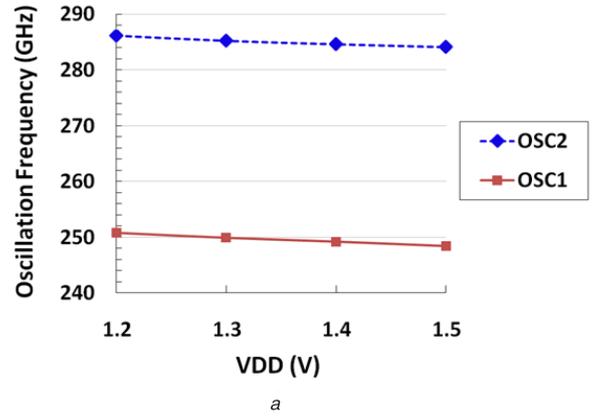


**Fig. 9** Testing setups  
(a) Power,  
(b) Frequency measurement,  
(c) Real (power) setup

0.7% frequency tuning range versus VDD with a ranging frequency from 284 to 286 GHz. Single-tone measurements at VDD = 1.5 V are given in Figs. 10b and c for OSC1 and OSC2, respectively.

Power measurements are corrected with the different setup losses summarised in Table 1. The resulting output power figures after losses extraction are depicted for both oscillators in Figs. 11a and b as a function of VDD. OSC1 maximum output power reached  $-8.1$  dBm at VDD = 1.5 V. OSC2 showed  $-14.8$  dBm output power under 1.5 V. DC power consumptions for both oscillators versus VDD are also given in Figs. 11a and b with maximum 76 and 80 mW for OSC1 and OSC2, respectively. Output power deviations between simulated and measured results could also be explained by the MOST modelling as part of the combined wires resistance was based on extrapolating the Calibre-found resistance to include the skin effect. Other reasons could be due to some over or under estimated shared connections or inter-metal layers vias or deviation from transistor parameters post-fabrication.

Phase noise measurements are carried out using the same measurement setup with the N9030A signal analyser. OSC1 gives at 10 MHz offset around  $-104$  dBc/Hz and  $-106.8$  dBc/Hz under 1.2 and 1.5 V supply voltage, respectively. OSC2's phase noise is



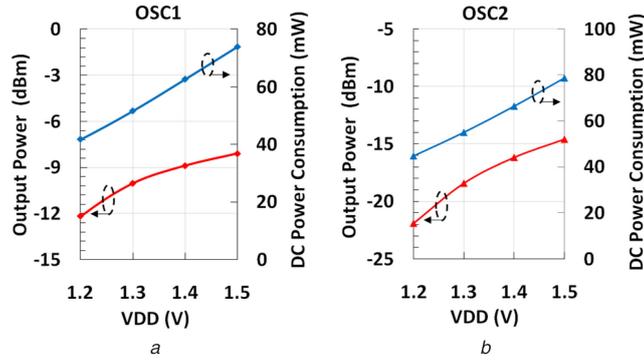
**Fig. 10** Measured oscillation frequency  
(a) Versus VDD for both oscillators,  
(b) OSC1 at VDD = 1.5 V,  
(c) OSC2 at VDD = 1.5 V

**Table 1** Setup power losses (in dB)

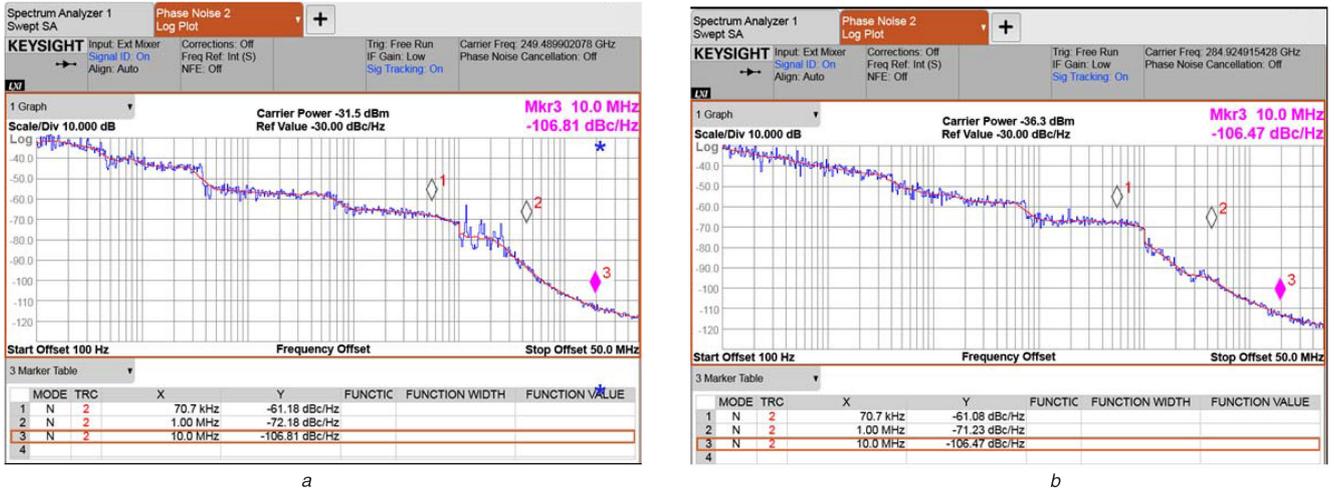
Circuit	Pad	I325_DL2CD (probe)	WR3/WR10 taper (3 cm)	WR10 WG (5 cm)
OSC1 (250 GHz)	0.5	3.0	0.35	0.375
OSC2 (285 GHz)	0.5	3.4	0.385	0.39

around  $-106$  dBc/Hz at 10 MHz offset for both 1.2 and 1.5 V VDD. Measured phase noise curves for OSC1 and OSC2 at VDD = 1.5 V are given in Figs. 12a and b, respectively. A plateau is present up to a few MHz offset due to the unlocked oscillation frequency. A slope between  $-30$  and  $-20$  dB/decade starts appearing from 10 MHz offset.

This work is compared with the state of the art of THz oscillators implemented in 65-nm CMOS technology in the 250–300 GHz range in Table 2. The presented single-stage oscillators in our case can be distinguished by the low DC power consumption but also by the lowest form factor. The proposed architecture can



**Fig. 11** Measured output power (after losses extraction) and DC power consumption versus  $V_{DD}$   
 (a) OSC1,  
 (b) OSC2



**Fig. 12** Measured phase noise at  $V_{DD} = 1.5$  V  
 (a) OSC1,  
 (b) OSC2

**Table 2** Published 250–300 GHz VCOs in 65-nm CMOS

Ref.	Freq., GHz	$P_{out}$ , dBm	$P_{DC}$ , mW	Phase noise, dBc/Hz	Area, $\mu\text{m} \times \mu\text{m}$
[7] <sup>a</sup>	239	-4.8	18.5	-98.43 to -110.9 @ 10 MHz	400 × 450
[9] <sup>b</sup>	256	4.1	227	-94 @ 1 MHz	670 × 650
[10] <sup>c</sup>	260	0.5	800	-78.3 @ 1 MHz	1500 × 1500
[11] <sup>d</sup>	288	-1.5	275	-87 @ 1 MHz	500 × 570
[13] <sup>e</sup>	298	0.9	235	-79 @ 1 MHz	290 × 316
[14] <sup>f</sup>	290	-1.2	325	-78 @ 1 MHz	600 × 600
this work (OSC1)	248	-9	76	-106.8 @ 10 MHz	200 × 250
this work (OSC2)	284	-15	80	-106.47 @ 10 MHz	200 × 200

<sup>a</sup>Single transformer-based push–push ( $2f_0$ ) VCO.

<sup>b</sup>Eight single-ended second-harmonic ( $2f_0$ ) VCOs.

<sup>c</sup>Eight coupled differential self-feeding second-harmonic ( $2f_0$ ) VCOs.

<sup>d</sup>Two mutually coupled single-ended triple-push ( $3f_0$ ) oscillator cores.

<sup>e</sup>Quad-core-coupled triple-push ( $3f_0$ ) VCOs.

<sup>f</sup>Four coupled differential fourth-harmonic ( $4f_0$ ) VCOs.

be evolved into a coupled-array source to increase the output power while keeping a low DC-power level as well as a small form factor advantage. This is an essential characteristic for large-scale THz systems integration in low-cost CMOS processes.

## 5 Conclusion

In this paper, a push–push oscillator architecture based on differential gate equalisation to enhance the oscillation frequency while achieving an ultra-compact design form factor has been presented. The frequency enhancement is derived as a function of

the equivalent RLC components of the oscillator's design parts. The design concept has been proven on two single-stage push–push oscillators implemented in standard digital 65-nm CMOS technology. Measured oscillation frequencies showed around 250 and 285 GHz with output power levels around -8.1 and -14.8 dBm, respectively. Phase noise levels are at -106.8 and -106 dBc/Hz at 10 MHz offset. The achieved design area below  $200 \times 250 \mu\text{m}^2$  represents a key feature for potential large-scale integrated coupled-oscillator arrays in low-cost and digital CMOS technologies.

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